

Effect of porosity on the Rippling motion of Herschel-Bulkley fluid in a non uniform tube

G. C. Sankad*, R. S. Metri and Asha Patil

Research Centre, Department of Mathematics
(Affiliated to Visvesvaraya Technological University, Belagavi)
B.L.D.E.A's V. P. Dr P.G.H CET, Vijayapur (586103) Karnataka, INDIA

-----ABSTRACT-----

The consequence of varying pressure drop on a Herschel-Bulkley fluid through porous tube has been explored. The governing equations of the model are analytically solved under long wavelength and small Reynolds number approximations. Results are obtained for the parameters: Darcy number, slip parameter, yield stress, axial radius on the pressure difference. The observations disclose that the building up of the porous thickening of the wall enhances the pressure. Also, rise in Darcy numbers reduces the pressure rise.

KEYWORDS: Peristalsis, Porous tube, Pressure drop, Slip parameter

I. INTRODUCTION

Peristaltic motion is the random movement of the muscles of the channel creating wave motions that drive the contents of the canal forward as they travel. In the digestive tract of human being, gastrointestinal tract, etc., smooth muscle tissues contract in succession, that generate a peristaltic wave propelling the food forward. Earthworms and some modern machinery use a similar mechanism for their locomotion.

Latham¹ was the principal investigator who studied the motion of a fluid in a peristaltic pump. The lubrication theory model wherein the model is considered to be inertia free and wall curvature is negligible, as discussed by Jaffrin and Shapiro². Mekheimer and Arabi³ examined the peristaltic flow inside a porous medium with MHD effect.

A non-Newtonian fluid does not follow Newton's law of viscosity. Numerous molten polymers and salt solutions are non-Newtonian fluids. Lately, the use of non-Newtonian fluids have become almost an essential need for the engineering, industrial systems and medical science. Therefore the research of diverse scientific mechanism with non-Newtonian fluids has also expanded. Medhavi⁴ studied rippling flow of non-Newtonian fluid. Eldabe and Abou-Zeid⁵ have analyzed mass and heat transfer of non-Newtonian fluid through a channel under rippling. Sankad and Patil⁶ have studied the non-Newtonian fluid flow inside a non-uniform conduit lined with porous peristaltic material.

The pores are naturally packed with a fluid (liquid or gas). Eldabe et al.⁷ has examined the non-Newtonian fluid flowing inside the horizontal channel under magnetic effect. Sankad and Dhange⁸ did analysis of the incompressible viscous fluid transport inside a peristaltic medium having pores, under chemical reactions and wall effects. Sankad and Nagathan⁹ investigated the heat transfer and slip impacts of MHD couple stress fluid through rippling motion under porosity.

In uniform tube the cross section of each stream of the tube remains unchanged and every particle progresses along its axis with constant speed. Srinivasacharya et al.¹⁰ considered micropolar fluid transport through rippling pipe. Vajravelu et al.¹¹ inspected the pumping of a Casson fluid in peristaltic tube having elasticity. Selvi et al.¹² analyzed peristaltic transfer of power-law fluid inside a flexible tube.

In a Herschel-Bulkley fluid model, the strain of the fluid is non-linearly associated to the stress. Vajravelu et al.¹³ analysed the rippling transportation of a Herschel-Bulkley fluid in a flexible tube. Rajashekhar Choudhari et al.¹⁴ discussed rippling motion of Herschel-Bulkley fluid in a flexible tube having porous walls under slip condition.

II. MATHEMATICAL FORMULATION

Herschel-Bulkley fluid flow inside a non uniform circular tube is considered under peristaltic motion with coordinate system (r, z, t) . The blood flow is modeled to be laminar, steady, incompressible, two-dimensional, axisymmetric and exhibiting peristaltic motion of Herschel-Bulkley fluid in an elastic tube of radius. The region between $r = 0$ and $r = r_0$ is called as plug flow region where $|\tau_{rz}| \leq \tau_0$. In the region between $r = r_0$ and $r = a(z, t)$, we have $|\tau_{rz}| \geq \tau_0$.

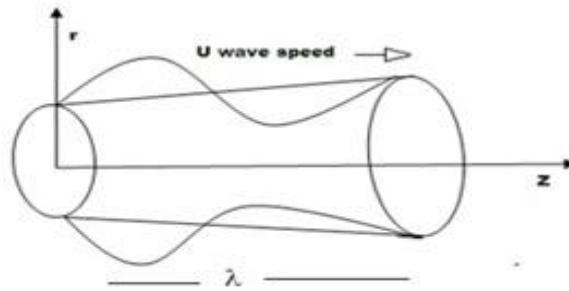


Fig. 1.1 Geometry of a symmetric peristaltic non-uniform tube with porous wall

The deformation of wall due to the transmission of rippling waves is represented as

$$H(z, t) = a_0 + b \sin \frac{2\pi}{\lambda} (z - ct), \quad (1)$$

where $a_0 = a + dz$ and a_0 : the mean radius of the tube; b : the amplitude of the wave; λ : is the wavelength; and c : the wave speed.

The constitutive equation for Herschel-Bulkley fluid is

$$\tau = \mu(\dot{\gamma})^n + \tau_0 \text{ for } \tau \geq \tau_0,$$

$$\dot{\gamma} = 0 \text{ for } \tau < \tau_0.$$

Governing equations

$$\frac{1}{r} \frac{\partial}{\partial r} (r \tau_{rz}) = - \frac{\partial p}{\partial z}, \quad (2)$$

$$\frac{\partial r}{\partial t} = 0, \quad (3)$$

where τ_{rz} is shear stress,

$$\tau_{rz} = \mu \left(- \frac{\partial u}{\partial r} \right)^n + \tau_0. \quad (4)$$

Here, n : the power law index; τ_0 : the yield stress of the tube.

The dimensionless quantities are

$$\bar{r} = \frac{r}{a_0}, \quad \bar{z} = \frac{z}{\lambda}, \quad \bar{u} = \frac{u}{U}, \quad \bar{q} = \frac{q}{\pi a_0^2 U}, \quad \bar{Q} = \frac{Q}{\pi a_0^2 U}, \quad \bar{a} = \frac{a}{a_0}, \quad \bar{p} = \frac{a_0^{n+1}}{\lambda \mu U^n} P,$$

$$\bar{\tau}_0 = \frac{\tau_0}{\mu (U/a_0)^n}, \quad \bar{\tau}_{rz} = \frac{\tau_{rz}}{\mu (U/a_0)^n}, \quad \bar{r}_0 = \frac{r_0}{a_0}.$$

The Eq. (2) is solved using the boundary conditions:

$$\psi = 0 \text{ at } r = 0, \quad (5)$$

$$\psi_{rr} = 0 \text{ at } r = a, \quad (6)$$

$$\tau_{rz} = 0 \text{ at } r = a, \quad (7)$$

$$u = -\frac{\sqrt{Da}}{\alpha} \frac{\partial u}{\partial r} - 1 \quad \text{at } r = a$$

$$= h(z) - \epsilon. \quad (8)$$

Here, u is velocity, Da : the Darcy number; α : the slip parameter and ϵ : the porous thickening of the wall.

III. SOLUTION OF THE PROBLEM

Using $p = -\frac{\partial p}{\partial z}$ in Eq. (2) we get $\frac{1}{r} \frac{\partial}{\partial r} (r\tau_{rz}) = p$.

Using Eqs. (5-8), Eqs. (2) and (4) are solved for the velocity field:

$$u = \left(\frac{p}{2}\right)^m \left[\frac{1}{m+1} ((k)^{m+1} - (r-r_0)^{m+1}) \right] + \frac{\sqrt{Da}}{\alpha} (k)^m - 1, \quad (9)$$

where $k = h - \epsilon - r_0$, $m = \frac{1}{n}$.

Using $\frac{\partial u}{\partial r} = 0$ at $r = r_0$, the superior limit of the plug flow region is given by $r_0 = \frac{2\tau_0}{p}$ and also using the condition $\tau_{rz} = 0$ at $r = 0$ we obtain $p = \frac{2\tau_0}{r_0}$.

Hence

$$\frac{r_0}{a} = \frac{\tau_0}{\tau_a} = \tau, \quad 0 < \tau < 1. \quad (10)$$

Using (10) along with $r = r_0$ in Eq. (9), the plug flow velocity is

$$u_p = \left(\frac{p}{2}\right)^m \left\{ \left[\frac{1}{m+1} (k)^{m+1} \right] + \frac{\sqrt{Da}}{\alpha} (k)^m \right\} - 1 \quad \text{for } 0 \leq r \leq r_0. \quad (11)$$

Integrate Eqs. (9) and (11) w.r.t. r and $\psi_p = 0$ at $r = 0$, $\psi = \psi_p$ at $r = r_0$ we obtain the stream functions as

$$\psi = \left(\frac{p}{2}\right)^m \frac{1}{m+1} (k)^{m+1} r - \left(\frac{p}{2}\right)^m \frac{1}{m+1} \frac{(r-r_0)^{m+2}}{m+2} + \frac{\sqrt{Da}}{\alpha} (k)^m r - r, \quad (12)$$

$$\psi_p = \int u_p dr = \left(\frac{p}{2}\right)^m (k)^m \left[\frac{k}{m+1} + \frac{\sqrt{Da}}{\alpha} \right] r - r. \quad (13)$$

The volume flux q :

$$q = \int_0^{r_0} u_p r dr + \int_{r_0}^a u r dr,$$

$$q = \left(\frac{p}{2}\right)^m \left\{ \frac{a^2}{2} (k)^m \left[\frac{k}{m+1} + \frac{\sqrt{Da}}{\alpha} \right] - \frac{(a-r_0)^{m+2}}{(m+1)(m+2)} \left[a - \frac{(a-r_0)}{(m+3)} \right] \right\} - \frac{(a^2 - r_0^2)}{2}. \quad (14)$$

From Eq. (14) will get

$$p = -\frac{\partial p}{\partial z}$$

$$= 2 \left[\frac{(2q + a^2 - r_0^2)s}{a^2(k)^{m+1}(m+2)(m+3) + \frac{\sqrt{Da}}{\alpha} a^2(k)^m s - 2a(a-r_0)^{m+2}(m+3) - 2(a-r_0)^{m+3}} \right]^{1/m},$$

where $s = (m+1)(m+2)(m+3)$.

$Q(z, t)$

$$= \int_0^H u(z, r, t) dr. \quad (16)$$

\bar{Q}

$$= q + 1. \quad (17)$$

Eq. (15) is integrating in respect of z over one wavelength to obtain the pressure drop

$$\begin{aligned} \Delta p &= \int_0^1 \left(-\frac{\partial p}{\partial z}\right) dz \\ &= \int_0^1 p dz. \end{aligned} \tag{18}$$

IV. RESULTS AND DISCUSSIONS

Herschel-Bulkley fluid having non-zero yield stress is examined to know the blood transport pattern in a non-uniform tube having elasticity. The yield stress considered, accompanied by the power law index, helps in exposing the shear thinning nature and thus inferring the blood flow characteristics.

The characteristics of the non uniform tube are observed by the graph of flow rate against pressure drop (Eq.18) using Mathematica software.

Pressure difference as plotted against \bar{Q} for varying values of Darcy number is represented in Fig. (1.2). decrease in the pressure drop Δp is observed with increasing Darcy number Da in pumping region and the opposite outcome is seen in the region of co-pumping.

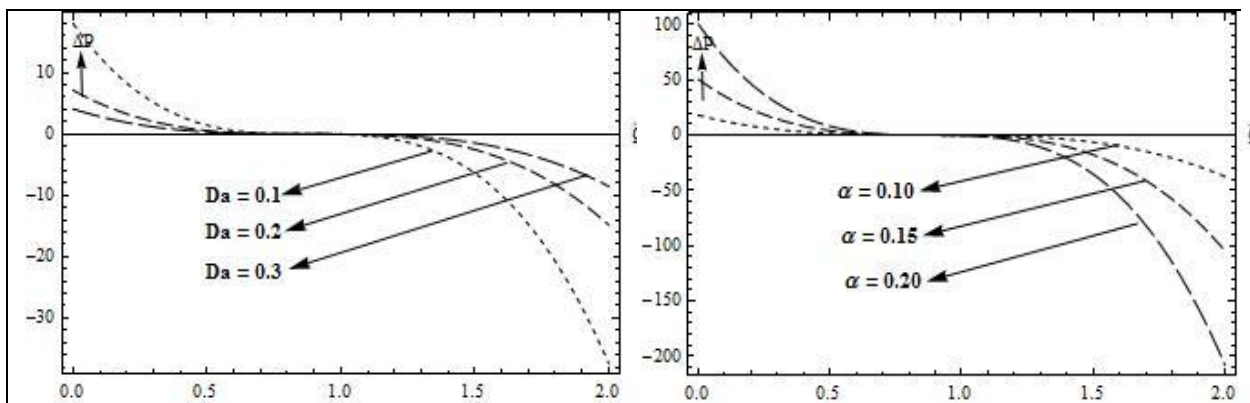


Fig.1.2.Outcomes of Darcy parameter.

Fig.1.3.Outcomes of slip parameter.

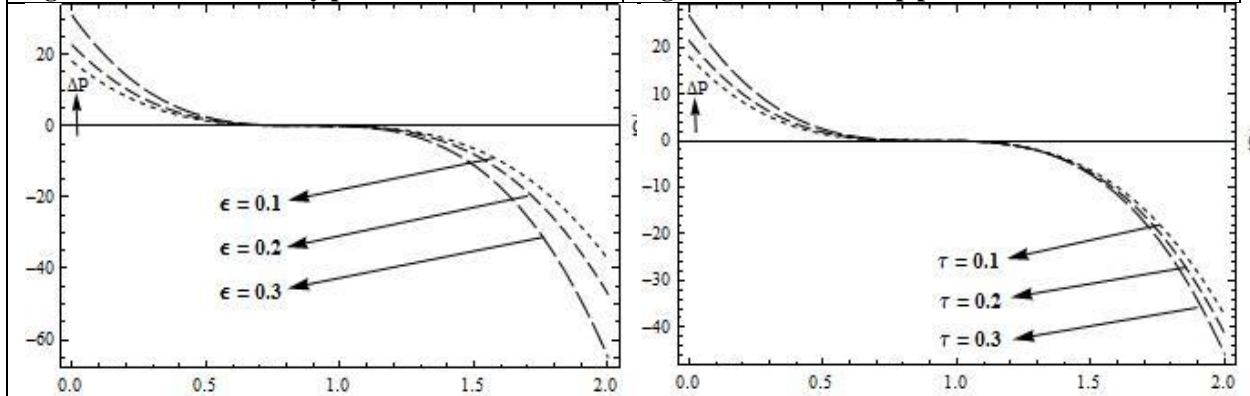


Fig.1.4.Outcomes of porous thickening of the wall.

Fig.1.5.Outcomes of yield stress parameter.

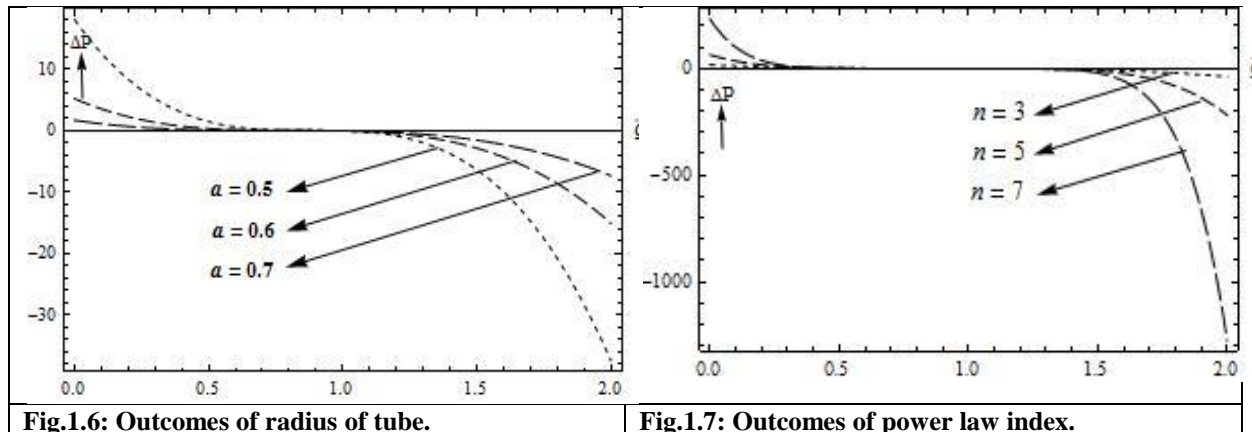


Fig.1.6: Outcomes of radius of tube.

Fig.1.7: Outcomes of power law index.

From Fig. (1.3) it is inferred that pressure drop Δp increases with increasing slip parameter α in pumping expanse and the contradictory result is seen in the co-pumping expanse.

It is seen from Fig. (1.4) that as the porous wall thickens the pressure drop Δp also increases in pumping expanse while effect is reversed in the co-pumping expanse.

The pressure difference variation with time averaged flow rate is considered with variations in τ as in Fig. (1.5). Increase in the pressure drop Δp is seen with increasing yield stress τ in the region of pumping while effect is contradictory in the co-pumping region. Observations infer that the peristaltic action on the tube wall, pumps with a larger pressure, than a power law index, because of the existence of plug flow region in Herschel-Bulkley fluid.

Fig. (1.6) depicts decrease in the pressure drop Δp with increasing radius of the tube in pumping expanse and effect is opposite in the co-pumping expanse.

Increase in the power law index n increases the pressure rise Δp in pumping expanse and the opposite outcome is seen in the co-pumping expanse as depicted in Fig. (1.7).

V. CONCLUSION

The Herschel-Bulkley fluid model is analyzed for the transport induced by sinusoidal peristaltic waves with small Reynolds number. Pressure drop increases with the power law index, slip parameter, yield stress and porous thickening of the wall but decreases with radius of the tube and Darcy number.

The above analysis provides acceptable results that signify some of the natural phenomena, mainly the flow of blood in arteries which can be processed and handled in case of dysfunction. The Herschel-Bulkley fluid is more emphasized as, blood behaves similar to Herschel-Bulkley fluid rather than power law and Bingham fluids, thus making it appropriate in the analysis of blood and other physiological fluid flows stimulated by peristalsis.

REFERENCE

- [1]. T. W. Latham, "Fluid motion in a peristaltic pump", MS Thesis, Massachusetts Institute of Technology, Cambridge (1966).
- [2]. M. Y. Jaffrin and A. H. Shapiro, "Peristaltic pumping", Ann. Rev. Fluid Dynamics 3, 13-36, (1971).
- [3]. K. S. Mekheimer, and T. H. Al-Arabi, "Nonlinear peristaltic transport of MHD flow through a porous medium", Int J Math Math Sci. (26):1663-1682, (2003).
- [4]. Amit Medhavi, "Peristaltic pumping of a non-Newtonian fluid", Applications and Applied Mathematics: An International Journal (AAM) Vol. 3, Issue 1, pp. 137-148, (2008).
- [5]. N.T. Eldabe, and M.Y. Abou-zeid, "The wall properties effect on peristaltic transport of micro polar non-Newtonian fluid with heat and mass transfer", Math Prob Eng, pp. 1-40. ID 898062, (2010).
- [6]. G.C. Sankad and Asha Patil, "Peristaltic flow of Herschel-Bulkley fluid in a non-uniform channel with porous lining", Procedia Engineering 127, 686-693, (2015).
- [7]. N. T. El-Dabe, G. Ismail, and F. O. "Darwesh, Peristaltic transport of a magneto non-Newtonian fluid through a porous medium in a horizontal finite channel", IOSR Journal of Mathematics (IOSR-JM) e-ISSN: 2278-5728, Volume 8, Issue 6, PP 32-39, (2013).
- [8]. Gurunath Sankad and Mallinath Dhange, "Peristaltic pumping of an incompressible viscous fluid in a porous medium with wall effects and chemical reactions", Alexandria engineering journal 55, 2015-2021, (2016).
- [9]. G.C. Sankad and P.S. Nagathan, "Transport of MHD couple stress fluid through peristalsis in a porous medium under the influence of heat transfer and slip effects", Int. J. of Applied Mechanics and Engineering, vol.22, no.2, pp.403-414, (2017).
- [10]. D. Srinivasacharya, M. Mishra, and A. R. Rao, "Peristaltic pumping of a micro polar fluid in a tube", Acta Mech. 161, 165-178 (2003).

- [11]. K. Vajravelu, S. Sreenadh, P. Devaki and K. V. Prasad, "Peristaltic pumping of a Casson fluid in an elastic tube", *Journal of Applied Fluid Mechanics*, Vol. 9, No. 4, pp. 1897-1905, (2016).
- [12]. C. K. Selvi, A. N. S. Srinivas and S. Sreenadh, "Peristaltic transport of a power-law fluid in an elastic tube", *Journal of Taibah University for Science*, 12:5, 687-698, (2018).
- [13]. K. Vajravelu, S. Sreenadh, P. Devaki, and Prasad K.V, "Peristaltic transport of a Herschel-Bulkley fluid in an elastic tube", *Heat Transfer-Asian Research*, 44, 585-598, (2015).
- [14]. Rajashekhar Choudhari, Manjunatha Gudekote, Hanumesh Vaidya, and Kerehalli Vinayaka Prasad, "Peristaltic flow of Herschel-Bulkley fluid in an elastic tube with slip at porous walls", *Journal of Advanced Research in Fluid Mechanics and Thermal Sciences* 52, Issue 1, 63-75, (2018).